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SIMULATION OF THE LOAD—UNLOAD PATHS EXPERIENCED BY ROCK IN THE VICINITY OF BURIED EXPLOSIONS

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1 December 1977

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TABLE OF CONTENTS

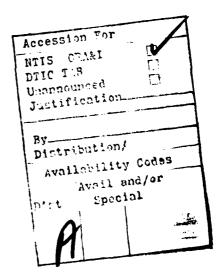
		<u>Page</u>
List	of Illustrations	1
List	of Tables	4
Intro	oduction	5
Stre	ss Path Determination from Finite-Difference Solutions	6
	Calculational Results	7
Stat [*]	ic Experimental Simulation of Load-Unload Paths	20
	Results	23
	ussion and Conclusions	35
_		37
Apper		
	General Relationships and Finite-Difference Calculations	37
	Analytical Determination of Elastic Stress and Strain Paths for a Spherical Explosion	42
Apper	ndix II	47
	Specimen Preparation	47
	Stress and Strain Determination	47
	Testing Procedures	48
	Data Acquisition and Analysis	48
	LIST OF ILLUSTRATIONS	
Figur Numbe		<u>Page</u>
la	Strain paths and stress paths at R = $2R_0$ cylindrical wave propagation in Mixed Company sandstone	9
16	Same as 1a, but with $R = 3R_0 \dots \dots \dots \dots$	10
1c	Same as 1a, but with $R = 5R_0 \dots \dots \dots$	11
2 a	Strain paths and stress paths at $R = 1.5R_0$ for spherical wave propagation in Mixed Company sandstone	12

LIST OF ILLUSTRATIONS (Cont)

Number		Pag
2b	Same as 2a, but with $R=2R_0$. Note changes in vertical and horizontal scales	13
2c	Same as 2a, but with R = $3R_0$. Note changes in vertical and horizontal scales	14
2d	Same as 2a, but with R = $4R_0$. Note changes in vertical and horizontal scales	15
2e	Same as 2a, but with R= $5R_0$. Note changes in vertical and horizontal scales	16
3	Strain paths and stress paths at various positions for cylindrical wave propagation in Mixed Company sandstone	. 17
4a	Comparison of the theoretical (calculated) strain path to the experimental strain path to be followed during testing of path I ($1/\alpha$ = 0.1 msec)	21
4b	Comparison of the theoretical (calculated) strain path to the experimental strain path to be followed during testing of path II $(1/\alpha = 1.0 \text{ msec})$	21
4c	Comparison of the theoretical (calculated) strain path to the experimental strain path to be followed during testing of path III ($1/\alpha$ = 10 msec)	22
5 a	Strain path followed during uniaxial-strain loading and constant-axial and uniaxial-strain unloading	24
5b	Stress path followed during uniaxial-strain loading and constant-axial and uniaxial-strain unloading	24
6a	Strain path followed during uniaxial-strain loading and constant-axial-strain unloading	25
6b	Stress path followed during uniaxial-strain loading and constant-axial-strain unloading	25
7a	Strain path followed during uniaxial-strain loading and constant-volume-strain unloading	26
7b	Stress path followed during uniaxial-strain-loading and constant-volume-strain unloading	26
8	Comparison of strain and stress paths determined numerically and analytically for spherical wave propagation in an elastic medium	46
9	Pressure vessel schematic showing the sample and	49

LIST OF ILLUSTRATIONS (Cont)

Figure Number		Page
9a	Stress path followed during strain path III testing	50
9b	Strain path followed during path I testing	51
9c	Stress path followed during strain path II testing	52
9d	Strain path followed during path II testing	53
9e	Stress path followed during strain path I testing	54
9f	Strain path followed during path III testing	55



LIST OF TABLES

Number	Description	Page
Ia	1236 Test Results	27
Ib	1239 Test Results	28
IIa	1241 Test Results	29
IIb	1257 Test Results	30
IIc	1285 Test Results	31
IIIa	1269 Test Results	32
IIIb	1270 Test Results	33
IIIc	1284 Test Results	34

INTRODUCTION

Common testing procedures for the laboratory measurement of material properties for use in ground motion calculations have generally consisted of standard hydrostatic, uniaxial-strain and triaxial tests. It has recently been recognized that these paths are not necessarily the ones that are followed in actual field applications, i.e., conventional and nuclear explosions in the earth. Since difficulty is often experienced in developing accurate constitutive models that are valid for a wide range of loading conditions, it seems important to follow, as closely as possible, the stress paths (or strain paths) that are experienced by material elements in actual field conditions. Furthermore, since measurement techniques do not yet allow the field determination of these stress paths (or strain paths), one must rely on numerical calculations and an initial best estimation of the material constitutive properties. In this report we present the results of one-dimensional numerical finite-difference calculations for cylindrical and spherical wave propagation, which define the stress and strain paths followed by material elements at varying distances from cylindrical and spherical explosive sources in the earth. The purpose of these calculations is to define laboratory tests best suited for the definition of material constitutive behavior in the analysis of CIST (Cylindrical In Situ Tests) and other subsurface explosive events. On the basis of these calculational results, static laboratory tests are conducted which represent strain paths experienced by material elements in the vicinity of cylindrical and spherical explosions in an infinite medium. The material tested in the experimental program is Kayenta sandstone.

STRESS PATH DETERMINATION FROM FINITE-DIFFERENCE SOLUTIONS

The quantities which are obtained from the finite-difference solution are σ_i and ε_i as functions of time at various distances from the explosive source. For purposes of definite laboratory tests, it is useful to express the output of these calculations in terms of the load L = σ_a - ρ_c and ρ_c in the triaxial test configuration. Here σ_a is the axial stress and ρ_c is the confining fluid pressure. It is also more convenient to deal with axial and transverse strain components (ε_a and ε_t) in the triaxial test rather than ε_i defined in the finite-difference solution. In the case of spherical flow, one would simply make the identification that L = σ_1 - σ_3 , ρ_c = σ_3 , ε_a = ε_1 , and ε_t = ε_3 . For cylindrical flow the identification is slightly more complicated.

In general, let us assume that we have values of stress and strain invariants defined by

$$\tau(t) = [(\sigma_1 - \sigma_2)^2 + (\sigma_2 - \sigma_3)^2 + (\sigma_3 - \sigma_1)^2]^{1/2} / \sqrt{6} , \qquad (1)$$

$$p(t) = (\sigma_1 + \sigma_2 + \sigma_2)/3$$
 , (2)

$$\varepsilon_{\nu}(t) = \varepsilon_1 + \varepsilon_2 + \varepsilon_3$$
 (3)

$$\varepsilon_{d}(t) = [(\varepsilon_{1} - \varepsilon_{2})^{2} + (\varepsilon_{2} - \varepsilon_{3})^{2} + (\varepsilon_{3} - \varepsilon_{1})^{2}]^{1/2} / \sqrt{6}$$
, (4)

as functions of time at a fixed spatial position as provided by the finite-difference calculation. If the material constitutive behavior involves only first and second invariants of the stress and strain tensors, the quantities defined by Eqs. (1) - (4) can also be written in the following terms for the purpose of defining laboratory test paths:

$$\tau(t) = (\sigma_a - p_c)/\sqrt{3} , \qquad (5)$$

$$p(t) = (\sigma_a + 2p_c)/3$$
 , (6)

$$\varepsilon_{v}(t) = \varepsilon_{a} + 2\varepsilon_{t}$$
 (7)

$$\varepsilon_{d}(t) = (\varepsilon_{a} - \varepsilon_{t})/\sqrt{3}$$
 , (8)

and hence laboratory stress and strain paths become in parametric form (t as the parameter):

$$L = \sqrt{3} \tau(t) , \qquad (9)$$

$$p_{C} = p(t) - \tau(t)/\sqrt{3} \qquad , \tag{10}$$

$$\varepsilon_{a} = \varepsilon_{v}(t)/3 + 2\varepsilon_{d}(t)/\sqrt{3}$$
 , (11)

$$\epsilon_t = \epsilon_v(t)/3 - \epsilon_d(t)/\sqrt{3}$$
 (12)

Calculational Results

Stress (and strain) paths for cylindrical and spherical wave propagation have been calculated with the use of elastic-plastic constitutive descriptions presented in the Appendix. The material parameters are chosen to be representative of Mixed Company (Kayenta) sandstone. In all cases a radial stress given by

$$\sigma_{r} = \rho_{o} e^{-\alpha t} \tag{13}$$

is applied at the interior cavity surface of radius R_0 = 1 m. The peak radial stress, p_0 , is taken to be 10 kbar and the decay constant, $1/\alpha$, takes on values of 0.1 msec, 1.0 msec and 10 msec. All results are presented in

terms of ϵ_a vs. ϵ_t (axial strain vs. transverse strain) and L/ μ vs. p_c/K (load/shear-modulus vs. confining-fluid-pressure/bulk-modulus), i.e., the quantities related directly to static triaxial laboratory tests.

Figures 1a, 1b and 1c show stress and strain paths at various distances from a cylindrical explosion. At the radial position R = $2R_0$ the stress path intersects the failure surface during loading and remains in contact during unloading. The corresponding strain path initially approximates conditions of uniaxial strain, but exhibits considerable transverse strain during the latter stages of deformation. At R = $3R_0$ it can be seen that the strain path is approximated by loading in uniaxial strain followed by unloading at constant axial strain, while at R = $5R_0$ the axial strain is seen to decrease during unloading. Of course, at much greater distances from the explosive source plane-wave conditions are achieved, and the load-unload path remains on the ϵ_{+} = 0 axis.

Figures 2a - 2e show similar behavior for spherical wave propagation. Figures 1 and 2 give an indication of how strain and stress paths depend on distance from the source. Another important consideration is that of pulse shape or pulse duration. This is controlled by the parameter α in Eq. (13). A number of calculations were performed for cylindrical geometry with $1/\alpha = 0.1$ msec, 1.0 msec and 10 msec. The peak radial stress p_0 remains the same in all calculations ($p_0 = 10$ kbar). The resulting stress and strain paths are shown in Figure 3 at radial positions $1.5R_0$, $2R_0$, $3R_0$, $4R_0$ and $5R_0$. One sees immediately that not only does the position have influence on stress and strain paths, but also that pulse duration has a significant effect. It will therefore be important to represent, as accurately as possible, the time history of the cavity stress due to the explosive source.

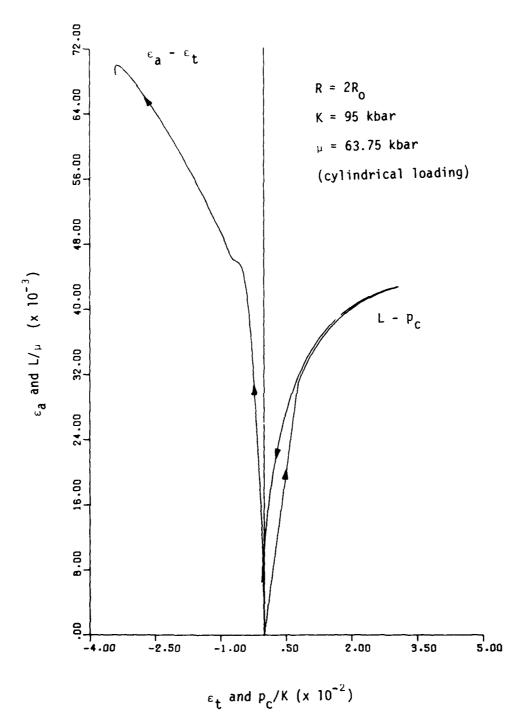


Figure 1a. Strain paths and stress paths at R = 2R cylindrical wave propagation in Mixed Company sandstone. A radial stress given by $\sigma_r = \rho_0 \exp(-\alpha t)$, with $(1/\alpha) = 1$ msec and $\rho_0 = 10$ kbar, is applied at $R_0 = 1$ m.

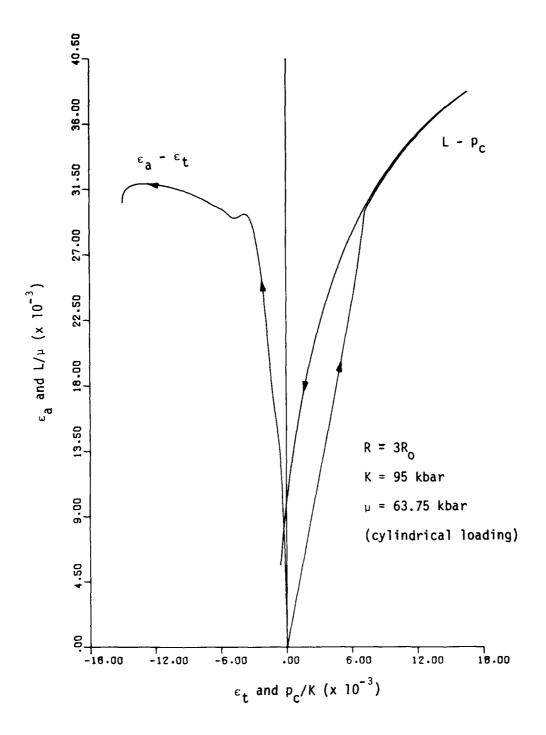


Figure 1b. Same as 1a, but with R = $3R_0$. Note changes in vertical and horizontal scales.

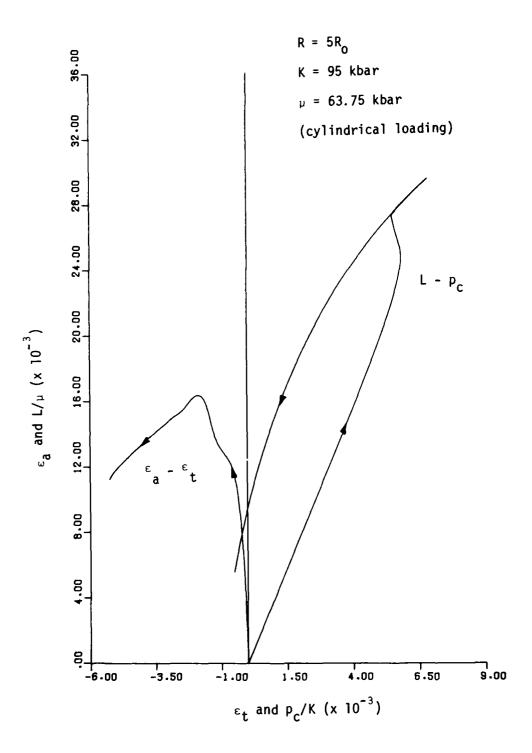


Figure 1c. Same as 1a, but with R = $5R_0$. Note changes in vertical and horizontal scales.

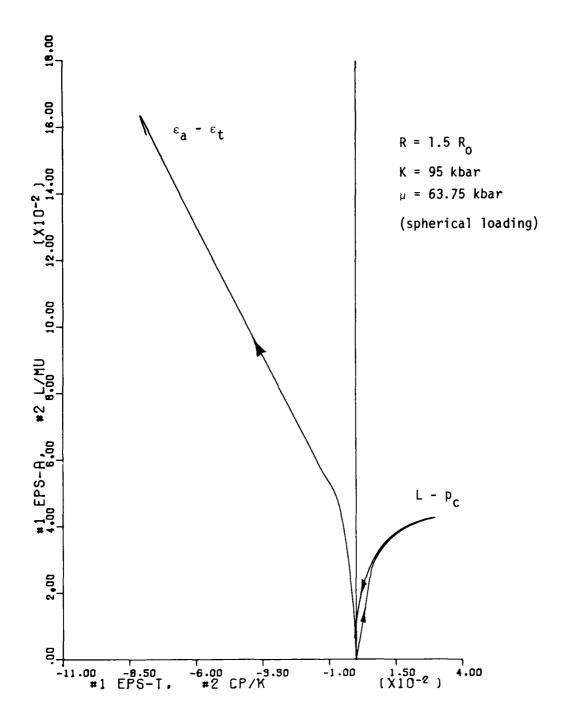


Figure 2a. Strain paths and stress paths at R = 1.5R $_{0}$ for spherical wave propagation in Mixed Company sandstone. A radial stress given by σ_{r} = p exp(- α t), with $1/\alpha$ = 1 msec and p = 10 kbar, is applied at R $_{0}^{0}$ = 1 m.

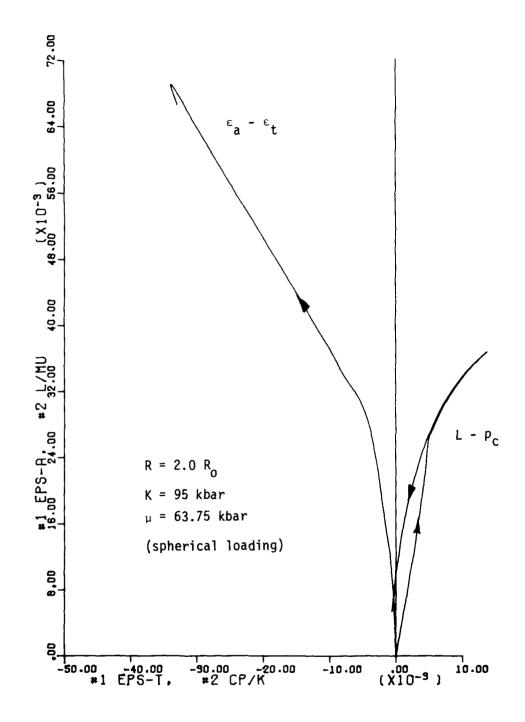


Figure 2b. Same as 2a, but with R = $2R_0$. Note changes in vertical and horizontal scales.

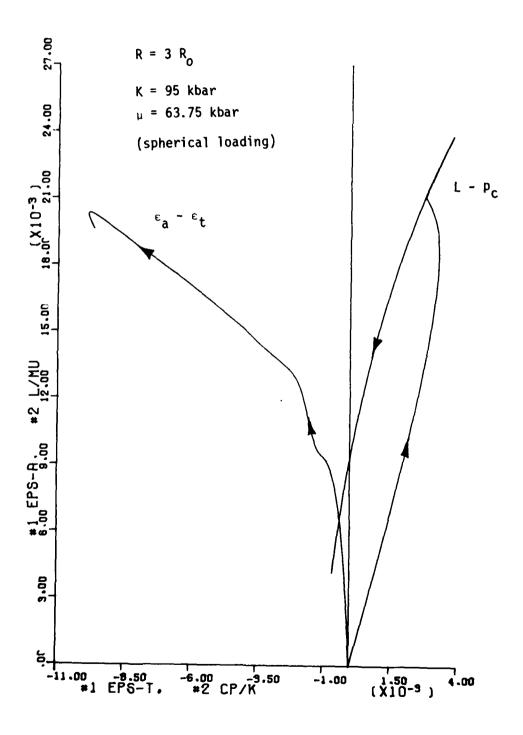


Figure 2c. Same as 2a, but with $R = 3R_o$. Note changes in vertical and horizontal scales.

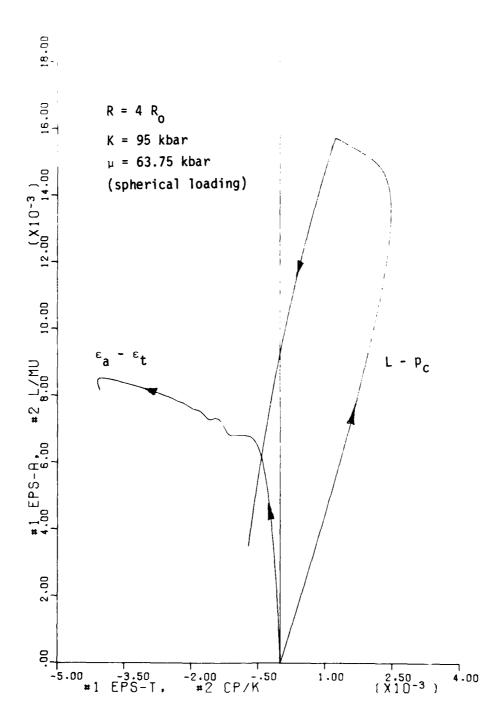


Figure 2d. Same as 2a, but with R = $4R_0$. Note changes in vertical and horizontal scales.

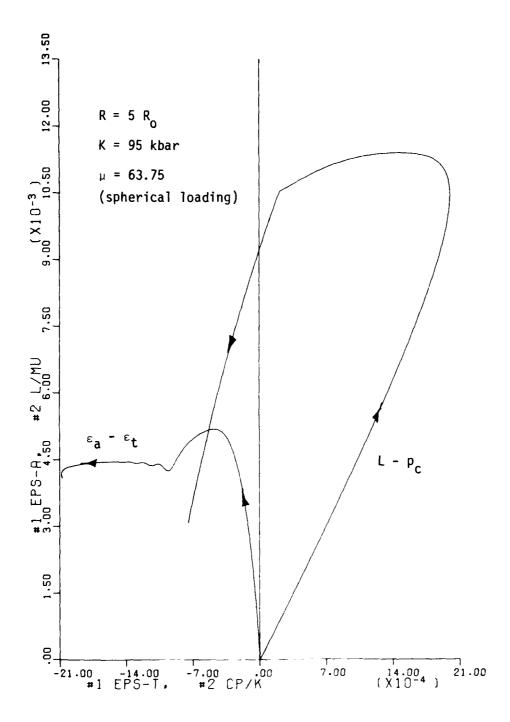


Figure 2e. Same as 2a, but with R = $5R_0$. Note changes in vertical and horizontal scales.

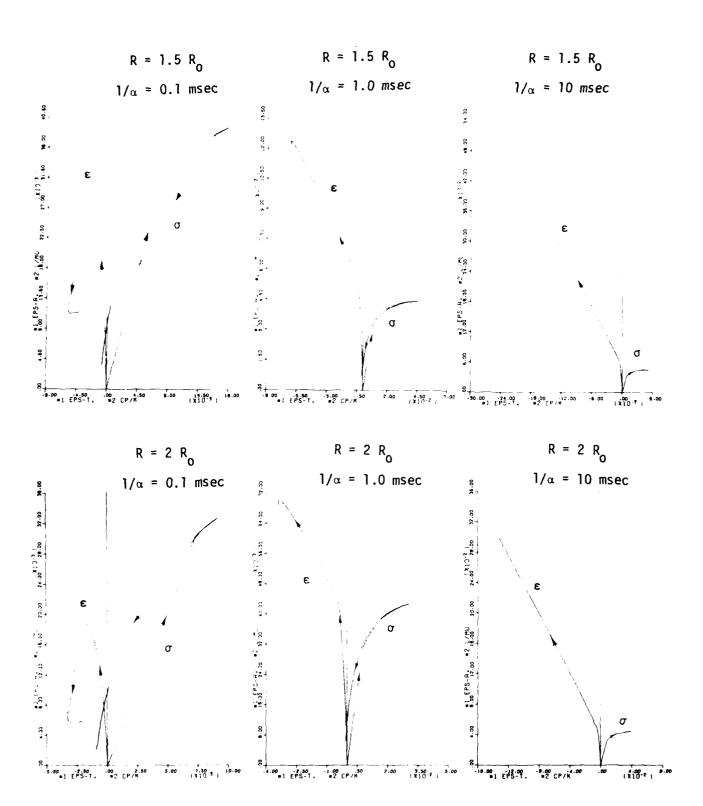


Figure 3. Strain paths and stress paths at various positions for cylindrical wave propagation in Mixed Company sandstone. A radial stress given by $\sigma_{\rm r}=p_0\,\exp(-\alpha t)$, with $p_0=10\,$ kbar and various values of $1/\alpha$, is applied at R = 1m. Note changes in the vertical and horizontal scales in each graph.

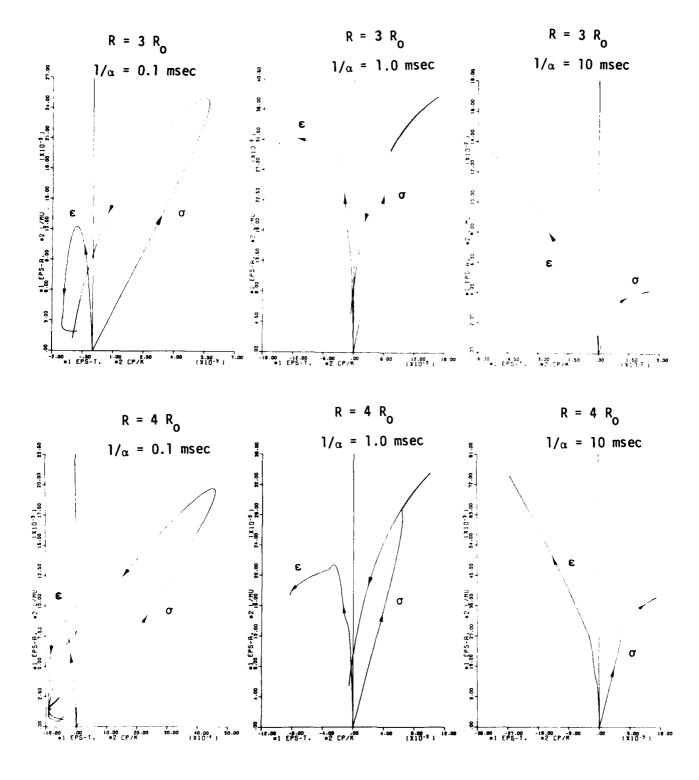


Figure 3. Continued.

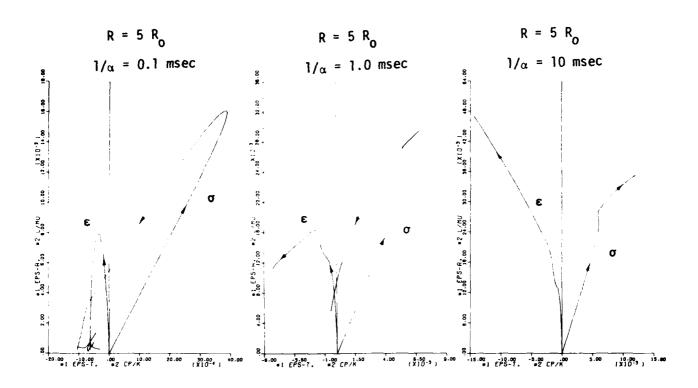


Figure 3. Continued.

STATIC EXPERIMENTAL SIMULATION OF LOAD-UNLOAD PATHS

Stress and strain paths were determined experimentally in the (L, p_c) and $(\epsilon_{a},\ \epsilon_{t})$ planes using the results suggested by various one-dimensional finite-difference solutions given previously. A detailed discussion of experimental techniques used in these tests is presented in Appendix II. The stress and strain paths considered here correspond approximately to those given in Figure 3 for $R = 3R_0$ and three separate decay constants $(1/\alpha = 0.1 \text{ msec}, 1.0 \text{ msec})$. Figures 4a, 4b and 4c show the three characteristic strain paths generated from the numerical solution and the strain paths to be followed in the static laboratory tests. The percent strains indicated here are used only to indicate orders of magnitude and are not the actual values achieved during testing. No attempt was made to follow the numerically determined strain paths exactly; they were used simply to indicate the qualitative nature of load-unload paths in the vicinity of buried explosions. Figure 4a shows the calculated and experimental paths corresponding to a decay time of $1/\alpha = 0.1$ msec; this consists essentially of uniaxial-strain loading and constant-axial-strain unloading followed by uniaxial-strain unloading. Figure 4b shows the theoretical path corresponding to $1/\alpha = 10$ msec in comparison to the experimentally followed path. The experimental strain path to be used consists of a uniaxial-strain loading and a constant-axial-strain unloading. Finally Figure 4c shows the theoretical path for $1/\alpha = 10$ msec in comparison to the experimentally followed path. The experimental strain path to be used consists of uniaxial-strain loading and constant-volume-strain unloading. Kayenta sandstone from the Mixed Company site was the material tested in this investigation.



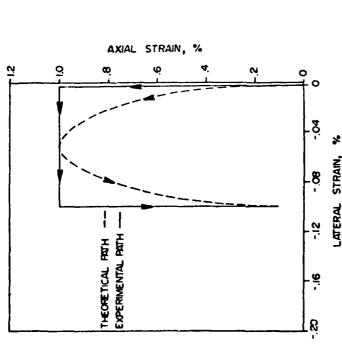


Figure 4a. Comparison of the theoretical (calculated) strain path to the experimental strain path to be followed during testing of path I $(1/\alpha \approx 0.1 \text{ msec})$. The experimental path shows a uniaxial-strain loading followed by a constant-axial and uniaxial-strain unloading. The percent strains indicated here are used only to indicate orders of magnitude and are not the actual values achieved during testing.

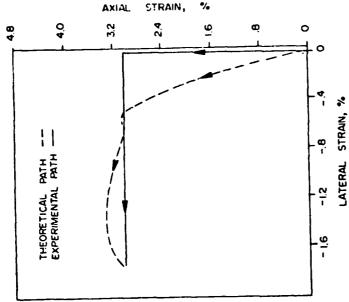


Figure 4b. Comparison of the theoretical (calculated) strain path to the experimental strain path to be followed during testing of path II ($1/\alpha$ = 1.0 msec). The experimental path shows a uniaxial-strain loading followed by a constant-axial-strain unloading. The percent strains indicated here are used to indicate orders of magnitude and are not the actual values achieved during testing.

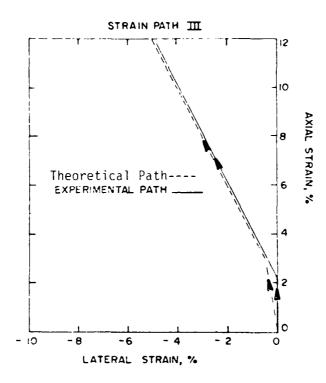


Figure 4c. Comparison of the theoretical (calculated) strain path to the experimental strain path to be followed during testing of path III ($1/\alpha = 10 \text{ msec}$). The experimental path shows a uniaxial-strain loading with a constant-volume-strain unloading. The percent strains indicated here are used only to indicate orders of magnitude and are not the actual values achieved during testing.

TEST RESULTS

The three strain paths, I, II and III, used in testing the Kayenta sandstone are shown in Figures 5a, 6a and 7a, respectively. Since all loading was conducted under uniaxial-strain conditions, a composite loading curve is shown for each path type. Individual unloading curves are shown for each test, departing from the composite loading curve at their respective maximum strains. The stress paths generated from the three strain paths are shown in Figures 5b, 6b and 7b. Composite loading curves are shown along with individual unloading curves. Included in each stress path figure is the triaxial failure envelope generated from this material. Tables I, II and III give computer listings for each test. Table Column 1 gives the data point while columns 2 through 8 give confining pressure (p_c) in kilobars, axial load (σ_a - p_c) in kilobars, axial strain (ϵ_a) in percent, the two transverse strains (ϵ_t and ϵ_t) in percent, volume strain (ϵ_a + ϵ_t + ϵ_t 1 in percent and mean stress [1/3(σ_a +2 p_c)] in kilobars. All plots were constructed from these tables.

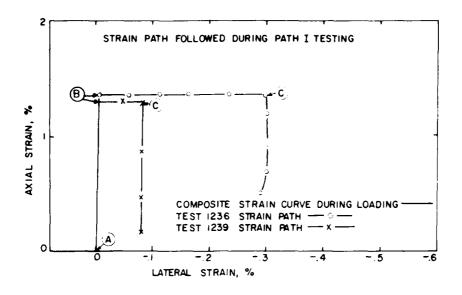
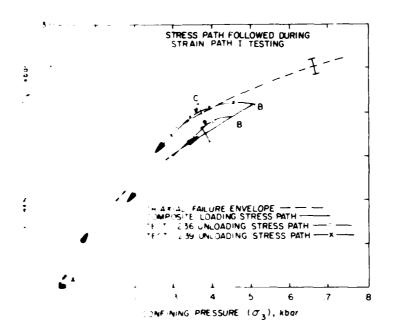


Figure 5a. Strain path followed during uniaxial-strain loading and constant-axial and uniaxial-strain unloading.



strain loading and constant-axial and uniaxial-strain unloading. The resulting stress path is a composite of four tests.

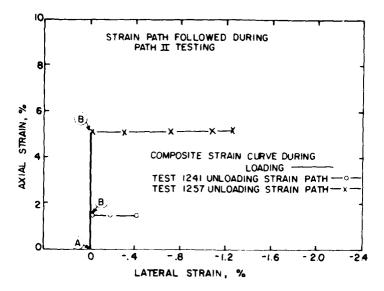


Figure 6a. Strain path followed during uniaxialstrain loading and constant-axialstrain unloading. The resulting stress path is a composite of four tests.

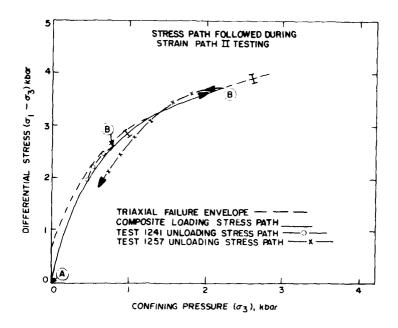


Figure 6b. Stress path followed during uniaxialstrain loading and constant-axialstrain unloading. The resulting stress path is a composite of four tests.

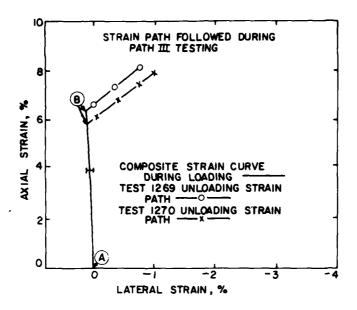


Figure 7a. Strain path followed during uniaxial-strain loading and constant-volume-strain unloading. The resulting stress path is a composite of three tests.

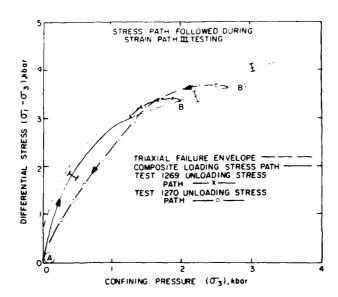


Figure 7b. Stress path followed during uniaxial-strain-loading and constant-volume-strain unloading. The resulting stress path is a composite of three tests.

TABLE Ia 1236 Test Results

Path Type I

CPRESS (KB)	LOPPO (KB)	EH (%)	E11 (%)	E16 (%)	VUL. STRMINGS	TENN SINESSINES
69	- 589743E-3	- 198966E-2	- 227923E-2	386529E-2	- 403623E-3	- 196581E-3
831515-2	121686	181541	282242E-1	- 249637E-1	164687	488504E-1
200948E-1	313309	389555	415941E-1	- 399487E-1	311286	124531
318745E-1	477203	40012	048219	- 584245E-1	397986	198942
5959156-1	728734	533219	571837E-1	- 613672E-1	529613	382583
789934E-1	86485	694486	572856E-1	- 867519	594191	367277
9146616-1	91214	637116	6190895-1	- 675545E-1	627434	395513
103246	949868	654895	632893E-1	- 660005E-1	652086	419868
121955	1 65335	708724	654594E-1	- 667998F-1	787375	473871
144821	1 16264	267879	641199E-1	- 71111F-1	768794	532367
177771	1 27770	924886	618403E-1	- 730928E-1	817547	597826
9461	9802	861657	654797F-1	- 045887R	658168	636886
342424	2020	100000	C04574E-1	CO24775-1	070000	22222
124517	2000	140000	A 10 10 10 10 10 10 10 10 10 10 10 10 10	1-35-1-200	200000	212213
20000	01000	30000	1 30000	1-34 (070)	10000	277.
25/8/5	1 48538	40.000	783873E-1	- E-50/14E-1	96167	731661
283466	1 60262	1 66219	6/0641E-1	- /80383E-1	1 86213	81/613
303501	1 6397	1 05002	674118E-1	- 694871E-1	1 64792	820068
11468	1 75746	1 16718	700496E-1	- 695277E-1	1 1077	913836
56,3095	1 88783	1 15244	681285E-1	- 700014E-1	1 15054	965702
192195	1 86739	1 19816	782766E-1	- 713021E-1	1 19712	1 01466
419912	40.40	1 23929	664379E-1	- 716842F-1	1 23398	1 6649
A55554	3/866	27.75	691769E-1	- 707975F-1	1 27116	1 19882
468417	2 62828	1 39454	719724F-1	- 686934F-1	1 20795	1 14468
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196424	62495	- 230715	- 258156E-1	- 161672	- 41c885	1 09974
N. 4000	65.239	5.63.3	- 428723E-1	- 177965	- 453154	1 07263
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160000	2 0113	- 238429	- 934331E-1	- 229259	- \$60692	1 03353
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. 5891	1 63442	- 291162	- 177727E-1	- 093062	- 401584	819899
257875	1 55954	- 317486	- 175714E-1	- 908614E-1	- 425546	776921
249493	1 49645	- 349849	- 158419E-1	- 899302E-1	- 455223	748269
223815	1 3761	- 483171	- 152192E-1	- 875477E-1	- 585497	682513
196.761	1 24882	- 46:275	- 198841E-1	- 925849E-1	- 575112	613065
178775	1 16985	56:93	- 216118E-1	- 918812E-1	- 616739	568723
156601	1 05114	- 56,114	- 203757E-1	- 913368E-1	- 674166	29638₹
139971	962498	- 694896	- 287744E-1	- 983536	- 71427	461884
124726	881336	- 654256	- 205561E-1	- 904281E-1	- 764484	418585
110175	989986	- 696613	- 217116E-1	- 696141	- 807656	379537
944.378E-1	702826	- 752216	- 193822E-1	- 866764E-1	- 857449	328513
824581E-1	55193	- 285669	- 240526E-1	- 934877E-1	- 902251	3028
658279E-1	554337	- 84975	- 277289E-1	- 914351E-1	- 967858	258687
519694E-1	458678	- 914621	- 261378E-1	- 852149E-1		284662
415755E-1	358952	- 995999	- 110911E-2	- 563872E-1	-1 65489	138339
, 1000						

* Axial strain rezeroed for constant-axial-strain unloading. ** Lateral strains rezeroed for uniaxial-strain unloading.

TABLE Ib 1239 Test Results Path Type I

U 56499E-3 - 266082E-2 - 227003E-2 36677E-1 27355E-1 365155 272692E-1 36677E-1 273892 366155 272692E-1 672615E-1 573415 473595 602774 672615E-1 573415 473595 60274 672615E-1 573415 473595 603774 115801 6736 445595 60408E-1 115801 10353 77249 264486 21466 10353 77247 264489 21466 12764 84466 225728E-1 21466 14114 95181 245489E-1 21466 141615 24548E-1 21426 15626 89466 22303E-1 21426 15626 89466 22303E-1 21426 15626 16645 21122E-1 21426 15626 89466 22303E-1 21426 15626 166446 226426E-1 21426 15626 166447 22		84968E-2 55204E-1 54189E-1 772188E-1	- 48119E-3 155626	- 194997E-3
19457E-1 25355E-1 156155 259828E-1 266677E-1 23382 306135 222935E-1 266677E-1 23382 415556 32695E-1 572615E-1 77213 417596 32656E-1 672615E-1 773195 36656E-1 3665E-1 672615E-1 773195 36656E-1 3666E-1 115801 8736 5266E-1 3666E-1 16431 16994 814643 22573E-1 16442 16994 74789 32666E-1 21426 17867 22782 32782 21427 1986 84468 22782 21426 17867 27786E-1 22782 21427 16845 227828 27123E-1 21427 16846 17666 27362E-1 21427 16846 17666 27362E-1 21427 16846 17666 27362E-1 21427 16846 17666 27362E-1 21428 16846 17666		55204E~1 54189E~1 83919E~1 70649E~1	155626	
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110841 00133 - 81412 - 520461E-1 - 88575E-1 560162 - 90572 - 506127E-1 - 596359E-1 400782 - 101419 - 506427E-1 - 506737E-1 - 506757E-1 - 5	•	12277	5.4316 -	416864
587573£-1 56015; - 9057; - 546127£-1 - 566339£-1 400782 - 101419 - 56627₹-1 - 5663377\$- 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337787\$ 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337₹-1 - 566337787\$ 5663378787\$ 5663378787\$ 56633787\$ 566	4	11411	- 981172	151659
596339E-1 400782 -1 01419 - 504237E-1 - 3	- 536127E-1 - 1	12354	-1 6.012	2,75465
1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	,	14639	-1 17692	193228
	- 505246E-1 - 1	112486	-1 25488	136326
18774F17 450468F11 -1 38785 -1 - 803087F11 - 820593E-1	•	320593E-1	-1 54786	226432E-1

* Axial strains rezeroed for constant-axial-strain unloading.

TABLE IIa 1241 Test Results Path Type II

z	CPRESS ++ B+	LORD (KB)	EA (%)	E11 (2)	ET2 (%)	VOL STRAIN	MERN STRESS(KB)
1	25	20468E-3	- 115851E-1	- 4442216-2	1727446-2	- 142995E-1	682267E-4
•	689171E-2	1998296-1	5,1612	- 544431E-2	- 294666E-1	497523	202194E-1
,,,	1722936-1	175272	617116	- 594697E-2	- 281611E-1	582801	749865E-1
1.	537694E-1	341891	661865	- 762541E-3	- 341489E-1	626785	147733
,	585795E-1	5.4288	767419	- 665779E-3	- 311719E-1	675358	236649
11	3-3897E-1	711335	775363	248558E-2	- 3252596-1	744608	328581
1.	31.8691	87.252.1	81791	- 865191E-1	- 3844735-1	77828	387889
ζ1	1716.2	3,115	878485	591776E-2	- 38,901E-1	865742	442216
·1	170225	1 06.54	(1.1.20 mg)	7-326690 -	- 379186E-1	16 36 † 3	524733
1,	195725	1 17792	415464	39,981E-2	1-35.0564 -	86969	588365
1,2		1 26715	3668	19459_E-2	- 495688E-1	916742	649122
	-61136	1 945	1 90,086	- 771885E-2	- 428818E-1	981819	72605
٠,٠	10816	F-00017 -	1 (0502)	·	- 1148255	1 64088	F07817
7	KT/	2.545	1 10847	446	1-48-44	1 67512	- C. C. C. C. C. C. C. C. C. C. C. C. C.
F: 7	64791	1 56.45	04017 T	5-3088069	1-3504075 -	30,50 1	32.6975
-4	W42014	1 76110	1 16.18	1-3+05T0 -	1-3865536-	1 10741	1 BO 399
1	472/846	1 81294	1 18447	ハーギザのハグオの エ	- 465508E-1	1 10651	1 65641
.7	117604	1 91.44	1 24501	5-3C7+528 -	- 4517546-1	1 19084	1 12575
:	54.	< 01071	1 V8×26	C-Jacoba C	- 6544	F. 25.57	1 51545
•	[* *	7.5138 3	1 . 1-4	7-379 871 -	1-27 53	1 5.7015	77.X1
7	75 10 0	15.14	# 175 T	- 1.00.6.16-2	1-30408-1	1 -011	1 -512.
,	C. 6. 616. 2. 5.	- 2045F	19.11	2-329C#58 -	- 11:2528-1	1,24564	1 34568
÷	6871184	F. 7. 10	1 19905	2-34004E-2	1-395.19.14	63ST T	1.42627
	1.650	5.11125	1 41169	- 101416E-1	- SUBSENSE-1	1 15664	1 47544
7. 7	17. 18.46.	A	※大学: · · · · · · · · · · · · · · · · · · ·	7-4-0044	- 14, 880.	5. TS. T	1 46.28
7		7177	1 402.0	174816971 -	1-30/2016-1	1 1 1 4 7	1 S1666.
) } {	74	るっナー	CHILDREST -	- 37,804E - 1	1 .7858	1 35.94.
9	₹	1084	17.7	アーナーのほん 一十	- 487 XESE-1	1 08952	1 Seves
	46134	2.41181	30,2 4 H	7-45 C C S -	- 411351E-1	7.364	1 666.1
•	†11 °	4.116	\$ 12×4 T	1 448995E-2	1-39000 7 1	1 40004	1 64174
	44/24	30707 ·	学士を行うして	2-400099 ×	1-47/34 4 -	「一旦の分する」に	1.6.201
ċ	γ. 2	82. 171	- 2010,084-1	- 18016.18-1	- 820571E-1	- 172904	1 5442
Ç	-c0114	22404 3	5-3154535 -	- 1514196-1	1 - 346×366 - 1	- 11101	1 62813
۲.	B1 (5 -	5 41119	- 2-81662E-1	- 485_e8t~1	335301 -	- 185179	1 5990.
j	4 4 4 4 4	12.00 m	0.000000000000000000000000000000000000	- 7915586-1	4.7.4	86898	1 56517
4		204-9-2	V-110+5.50 -	- 1878.6	F. 1. 1881 1	- 297161	1 5392
:	7.15		7,71816-2	7	7.11.11	- 4. F. C. Y.	7,447

* Axial strain rezeroed for constant-axial-strain unloading. ** Sample failed due to jacket leak.

TABLE IIb 1257 Test Results Path Type II

ź	CPRESS (KB)	LOAD (KB)	ER (%)	E11 (%)	ET2 (%)	VOL STRAINCES	MERN STRESS(KB)
-	0	- 148644E-4	- 869849E-2	- 611251E-2	- 624006E-2	- 210496E-1	- 495481E-5
1	179088F-1	081949	1354	- 158112E-1	- 341167E-2	116152	452251E-1
	365063E-1	24764	296826	- 012948	- 738807E-2	276431	119053
	723239E-1	482886	44911	- 838544E-2	- 825744E-2	432393	233286
7.	201129	1 19917	799277	- 089785	- 140653E-1	775318	567852
11	387264	1 38882	984475	- 573949E-2	- 01274	965814	770144
	497313	1 87199	1 3003	- 794519E-2	- 194789E-1	1 27252	1 12131
	546986	1 9486	1 35369	- 351501E-2	- 169373E-1	1 33296	1 19625
4	579969	2 64992	1 4176	- 514466E-2	- 188999E-1	1 39322	1 26328
1.9	680533	2 26515	1 56354	- 389695E-2	- 821617	1 53763	1 43558
<1	725994	2 33251	1 61083	- 172212E-2	- 203504E-1	1 5884	1 5035
, ,	.63189	2 48457	1 66431	- 261109E-2	- 247494E-1	1 6365	1 56471
۲,	789361	2 46736	1 76425	- 189321E-2	- 248209E-1	1 67798	1 61182
	365738	2 50936	1 75592	121275E-2	- 233032E-1	1 73345	1 66439
Ţ.	885735	2 63541	1 84435	632482E-3	- 261527E-1	1 81836	1 76427
	996459	2 68121	1 87655	- 178195E-2	- 269812E-1	1 84725	1 8002
	86736	2 71291	1 90657	- 211972E-2	- 273343E-1	1 87655	1 82729
J.	568845	\$157.64 \$157.64	1 90109	- 494927E-3	- 023347	1 90679	1 65215
	97.53.59	63/1+1 N	1 96086	610578E-2	- 234102E-1	1 94322	168 1
	1 02218	99068 7	2 842	256024E-2	- 262752E-1	2 8178	1 9734
41	1 6469.7	2 90255	2 69428	110712E-3	- 252145E-1	2 06866	2 61449
7	1 (4965)	2 95722	2 15461	251423E-2	- 222904E-1	2 13441	2 87564
4	1 1568	24460	2 31133	- 335846E-2	- 266806E-1	2 2806	2 1983
+	1 20678	0.11995	2 40735	- 205731E-2	- 240868E-1	2 38958	2 24676
7.7	1 2257	15545	2 5093€	- 328668E-2	- 233136E-1	2 48217	2 28752
1,	1 28.54	550LT 1	2.689.3	- 729194E-2	- 220813E-1	5 65956	2 33941
3,5	1 71836	17925	2 83887	- 721278E-2	- 181939E-1	2 81274	2 37811
£	1 75969	0.19787	3 94567	- 111841E-1	- 185468E-1	3 01504	2 42431
٠.٠	1 586, 7	2.55647	2 77984	- 1796528-1	- 110108E-1	3 74897	2 76535
£.Ç	1 70684	2 41846	4 19266	- 216131E-1	- 87089€E-2	4 97169	2 84366
1.5	1.93552	24562	4 48361	- 251871E-1	- 285185E-2	4 4528	1172
ú	2 17247	6,7459	4 8432	- 298487E-1	- 429919E-2	4 81241	44/C04 I
4	2 29921	3 73457	5 018	- 284185E-1	323877E-3	4 9885	3 54487
i	38738	2 79135	5 11999	- 2784136-1	257092E-2	5 69345	3 65115
7.4	2 38748	7.672	- 335449	- 509426E-1	- 291423E-1	- 415249	3 55673
	2.1714	75040	- 228219	- 116933	- 977041E-1	- 542037	3 38229
	· 9536.4	3 71152	- 15964	- 156404	- 141204	- 612352	3 26086
ď.	1 7.7.3.		- 291567	- 254155	- 240676	1 784345	2 96287
1:	1 45612	28461	- 279294	- 467839	- 392585	-1 07589	2 55699
*	989115	2 65863	- 244605	- 835432	- 22266	-1 84192	1 87553
10	724616	2 15215	- 21545	-1 38836	-1 12775	80.862 2-	1 43533
	é51911	1.80,04	- 16.892	- 117778	- 194421	- 475336	1 30626

* Axial strain rezeroed for constant-axial-strain unloading.

TABLE IIC* 1285 Test Results Path Type II

	CRAC SCARAGO	LOME - P.B.	1	£.1	E12 - 25	VOL STRRINGS	MENN STRESSINES
	œ	1 285825E-3	- 389069E -	- 663233	- 2687288-2	- 981865E-2	- 127942E-3
	347881E-2	233786E-1	493756E-1	- 611584E-2	- 858556E-3	423979E-1	921902E-2
	1-38896-1	255.5	58634	- 117936E-1	486112E-2	351877	111082
	55226E-1	627454	586816	- 118582E-1	13665E-2	576262	264875
	一日の一日の日の日の日の日の日の日の日の日の日の日の日の日の日の日の日の日の日	780215	687141	- 157566E-1	- 2673238-2	58282	53437
	677131	857.168	#368CC	- 12 4078-1	- 916211E-1	756683	369861
	122878	¥86647	847272	- 116234E-1	- 6518356-3	E34893	45176
	900191	1 18465	808606	- 789414E-2	694606E-2	932625	531817
	11000	1 19603	1 01532	- 120642E-1	- 421843E-2	99894	581727
2	3000 C	1.45.45	1 12042	- 612757	- 486564E-2	1 1156	683024
	(1) (1) (1) (1)	1,500,1	CC062 T	- 10×119E-1	- 1161856-1	1 26597	8166
	すい。まなず	97.00t	1 50811	- 1:1252E-1	- 725151E-2	1 48,44	1 90554
	400044	1 46788	1 61617	- 160811E-1	- 959154E-2	1 5846R	1 685.2
	たいめたたす	1 96145	1 65124	- 124684E-1	- 784195E-2	1 6386	1 15175
	のすのいすの	2 10477	1 78364	- 188918E-1	- 911693E-2	1 73555	1 24953
	616817	25015	1 90768	- 192815E-1	177829E-1	1 90525	1 36067
	11.	137 ·	2 60.667	- 195538E-1	- 851146-2	2, 660,44	1 46647
	\$ 13.	の意味が大い	48260 7	- 186077E-1	- 591604E-2	2 6/281	1 5-122
	528282	2.53682	2 2117	- 196729E-1	- 291889E-2	2 18759	1 63288
	846141	2 64091	2, 02879	- 2271725 -	- 548768E-2	2 29194	1 72044
	841688	2000	2 06269	- 196521E-1	- 977153E-3	2 34157	1 77139
	14444	1609C /	2.46725	- 1914881-1	- 121856E-2	2 44636	1 84475
	154446	A4400 0	\$1.75 V	- 6286 · 2	1-男子7:1	2.50217	1 88994
	1. 1.400	000,000	26.769	「「一般ののから」。	7-B800097 -	2 66832	1 34403
	1 8615	NA 2	2.81799	- 226248E-1	- 749888F-1	2 78795	1 98134
	1 6 6	2 88011	2 89016	- 249607E-1	- 143657E-2	2 863	1 99307
	1 68918	38865	12882	- 208663E-1	- 163973E-2	3 10561	2 64953
	たんの女子 エ	2 44881	9669 3	- 185413F-1	- 391683E-4	3 55034	2 17171
7	1 55	6.000	305 an	- 199897t-1	C-3241141 -	3 84476	2 22748
0	1.00	15011	16967 7	- 214228E-1	- 556425E-2	4 25879	61.79
	40000	2001	4 41419	- 227916E-1	- 848947E-	4 38153	2 41567
	7.75 1	5.5516	4 99419	- 207557E-1	- 112912E-1	4 96055	2 63127
	1 66171	3,35202	5 41/54	- 213594E~1	- 123665E-1	2 38143	0000
_	1 78744	3 42621	5 84727	- 170119E-1	- 139484E-1	5 8145	2 92951
	1 8646	1.60°7	6 66788	- 181019E~1	- 1055346-1	6 05427	. 02658
	TERRE T	371.6	6 47219	- 61784	- 123867E-1	6 44086	1894
	2 18464	39869	6.72226	- 199912E-1	- 743502E-2	6 69299	5 34752
		1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	44.00	1 GC 204E - 1	0-0644500	11,000	

* Showing only the uniaxial-strain loading.

TABLE IIIa 1269 Test Results Path Type III

z	(PPE++ +) B+	Lindo o Pers	I	671 +2:	E12	VUL STERIN X	VUL STRAIN X) MEAN STRESSYB)
	0	- 227e725 -	- 1077788-2	- 820014E-3	229861E-2	982016E-4	- 758773E-4
1 - 4	27675E-2	125848	257922	120078E-1	- 263284E-2	24732	447117E-1
3	621531E-2	174692	285465	225077E-1	585931E-2	313915	064446
	1312126-1	31689	405245	2180465-1	418729E-2	431344	119418
-1	27623HE-1	5,0053	554722	253035E-1	759655E-2	587887	204308
v	4:50.16-1	6, 36%	550555	256479E-1	1018956-1	699021	269869
11	5771896-1	サベニバチト	7.17. J	289758E-1	128151E-1	772547	321447
	7596.48E-1	をすった こめ	の すべ のの	285461E-1	250991E-1	856573	368414
15	966825E-1	920366	868731	258486F-1	366556E-1	941692	416141
1	114608	1 08569	いずかいの	110885E-1	980034E-2	53857	476533
7	144	1.18785	1 (15,7%)	148972E-1	192485E-1	1 69266	528951
	1.72.5.4.	1 .8601	1 150%	1:4::1E-1	1-3890477	1 1877	601016
	36661	1 4114	1.77.1	1十五十二十十二十二十二十二十二十二十二十二十二十二十二十二十二十二十二十二十二	2873196-1	1 26790	652140
٠,	はまれる	7.644.1	7.97	1454636-1	14529E-1	1 35421	2,100,15
	266586	4	1 +6711	1864788-1	416974E-1	1 46622	778016
	6164	1 58.6	40000	198122E-1	474544E-1	1 62284	828974
-	11.1	1 66.54	172451	016.61	W51400	1 29049	866350
	PARTY N	# 72	1 85.7 1	- B. S. C. 1966	1,601,612条-1		901807
	-100	+ -T 1	18:00	219770 -	10:45	J 98415	928645
	71017	1.65.21	+10th 1	1-30232F-1	12294	PCP281 3	369568
,	1. A. A. A. A.	1 7.00	1,000	11-日からかから -	161895	7 16241	1 65799
+1	-1,15	1 80045	サナナロ	: 74966.1-1	175090	56940	1 11532
7	55.55	*	1 60.00 Ac	1 - 470,000 -	198158	2 34721	1 17942
4	241175	; %	1.000	T - 4708/3005 -	6.117. 6	J 41871	1 22246
7	77	7.7 1	Columbia.	- 1000te.5	95,75,27	2, 512.92	1 27755
•		*********	144	F00001 -	# 122 7 2	2.58877	1 0020
7-	92.5	2 (6.15)	311.15	- 1106.4	267596	2 72903	1 41415
,	1 A	2421 2	1000	- 115,54	2712	2 81845	1 46549
4.	1 (4)	1.07	and and a contract of	74514	34527	3 56942	1 82604
• .		114.54	7.5° - 7	- 149681	7 197	1956	2 25036
7	1	1184	4 8445	- 117,64	#1112 7	5 17874	2.2551
7.	4	34146	東京大学	1 - 4 - Mar - 4 - 1	気のたいすす	5 20424	4.558
i	\$4.1. ×	5 28114	6 208es	- 164'87E-1	かんしゅうす	F. 74653	2 81326
ig' i	16155	10.	7 6 1 0 1 0 1 0 1 0 1 0 1 0 1 0 1 0 1 0 1 0	1 0000000	1 504007	465669	መ ነገር ነው። የመርቀ የመርቀ የመርቀ
an L 10	0.350		0100	90000	-1 74108	959807	16950
-	· · ·	7 3	1	1 77		30	17466
	7 7 7	177.5	چ پ	Sept. 24 1-	53563 T-	26198	11952
117	11.75.00	はすすす()	4 17.35	9861 7-	-1 9452	化价的量	2 7762
11	1 5510		5797 7	1962 1-	-1 95815	45635	2 62902
114	18 SO 1	. 883.5	4 184	-1 33318	1. 9.757	42945	2 23866
115	1 7 (8)	4152.7	14:17	6890 Tr	1.1000 Tr	41633	2 12768
115	1 +++1	2.0	7 4 7	-1 7565	- 6 19617(1)	425.78	1 8:858
117	113678	Street, 15	1 1 1	7.7	-2. British	2.833	1 57517
118	187854	1 86461	+ 18e	-1 72942	-2 04316	41342	1 32946
11.5	€2259	1.69065		-1 72468	-2 64935	37897	1 21/85
171	591145	1.5%.	4 147	-1 71918	-2 0565	36792	1 16591
14	721815	1 016, 4	\$ 555 E	-1 69415	77688 2-	1574	95,099

* Axial and lateral strains rezeroed for constant-volume unloading.

TABLE IIIb 1270 Test Results Path Type III

The second of the second secon

z	CERESS FE	LOHO + KB3	ER (2)	E11 (2)	ET2 .%,	VOL STRAIN(Z)	VOL STRRINCZ) MERN STRESSCKB)
-	2	1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	- 10 adde-2	1 8208148-1	2298616-2	814884E-4	~ 758713E-4
•	4.4.2.24	- サテクロ・・・	75.00	1.000 E	2801526-2	2139 05	719988E-2
		FOR PATE-1	44.300.7	141636E-2	2203376-2	265326	291677E-1
. 4	1-34646-1	1.551	000000	45.6980E-2	264551E-2	329244	940242E-1
, ,	15.598.7E-1	40Me 7	7000	S_8375E-2	228105E-2	3977.68	162156
٠.	1-4000	7.00	4856.	811500E-2	241429E-2	196264	26232
	1614.	7	P. C. P. P. P.	1十分の公司の第一1	1692786-1	5.2.389	499464
	7.27.77	10000	4	1.9.7.6-1	01156	861182	567165
•	1 44.W.	14.1	T1177 7 14	15590.6-1	2458746-1	1 01048	614719
141	73.	4111	1.1817	215469E-1	412426E-1	1 24491	789954
: =	+114.114	1 - 1 - 1	1 444 1	1 - 3+14+14 C	508888E-1	1.5776	1 4128E
	-1.1.	7		1-元 できます	8572846-1	2 56246	1 :1986
-	- 4.4	1-+++-1		47118, 1-1	1-344-10-	#12PF 2	1 252
. 1	;	4 41	7.	5 70925-1	1-4-10-5-	2 21917	1 42058
	î		14.	E-3414E-1	9, 214,26-1	1.00500	1 50564
	7.17	7.	7.1.7	107175	1-4-7-20	2,040,05	1 : 86.64
	•		7.7	1111.4	1-34-46-70	55,59	1,79954
. ~	1848	÷	7.	2000	81576.4 -1	46498	1 91 4
	7	7-7	ii T		I gradedie		L.Come of
			1		To Branch Code	14. S. #	2 21862
•••	7		4 . 14-1	140.56	11,501.9	5 1154.	46.00
		+		14.14	145151	5 476	2. Section 2
	+	• • • • • • • • • • • • • • • • • • • •			15414 14.	1000000	S 541.1
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* Axial strain rezeroed for constant-volume unloading. ** Could not maintain constant volume path beyond this point.

TABLE IIIC* 1284 Test Results Path Type III

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*	3, 00135	4 24656	7 9991	- 209629E-1	184262E-1	2 88873	4 41687

* Shows only the uniaxial-strain loading.

DISCUSSION AND CONCLUSIONS

An initial observation of the experimentally observed stresses indicates that there is little difference between the load-unload paths for strain path types II and III. Figures 6b and 7b suggest that if the yield condition is reached during uniaxial-strain loading with the stress paths following along the yield surface, then the unloading stress paths are similar in direction and magnitude for either constant-axial-strain or constant-volume-strain unloading. The numerial analysis solutions agree with the above observation in that regardless of the strain path, the stress path would follow along the yield surface during unloading (provided that yield was reached during uniaxial-strain loading). All of the experimentally observed stress paths show the unloading curve to go initially above and then cross through and go below the loading curve. The experimental unloading curves did not remain on or intersect (as in the case of strain path 2) the yield surface as illustrated by the numerical analysis.

Such variations in unloading material behavior may be modeled by including additional phenomena into the constitutive equations. Phenomena to be included in the equations would be permanent volume compaction and work-hardening of the shear failure envelope. The former effect will mainly influence the strain paths and the latter will change the stress paths, particularly in the unloading portion. It was experimentally determined that the material behaved nonlinearly during initial loading as compared to the linear model used in the numerical analysis. Such nonlinearities may be also handled by the aforementioned considerations. The observation that the unloading path lies below the loading path in stress space may be related to fracture and the resulting loss of cohesion, rather than ductile plastic flow, as assumed in the calculations.

Inclusion of pore pressure effects into the model would be of interest in future work. Both the calculations and laboratory strain-path tests should be performed under various saturation conditions. Much of the previous theoretical work, including the finite-difference computer code, already contains this capability; it has just not been exercised yet. Also of future interest would be some theoretical results for two-dimensional dynamic loading situations, expressed in terms of ε_a , ε_t , L and ρ_c . This could be done by calculating the following invariants as functions of time at a particular material element:

$$\tau(t) = \left\{ (1/6) \left[(\sigma_{11} - \sigma_{22})^2 + (\sigma_{22} - \sigma_{33})^2 + (\sigma_{33} - \sigma_{11})^2 \right] + \sigma_{12}^2 + \sigma_{13}^2 + \sigma_{23}^2 \right\}^{1/2}, \tag{14}$$

$$p(t) = (\sigma_{11} + \sigma_{22} + \sigma_{33})/3$$
 , (15)

$$\epsilon_{\mathbf{v}}(t) = \epsilon_{11} + \epsilon_{22} + \epsilon_{33}$$
 (16)

$$\varepsilon_{d}(t) = \left\{ (1/6)[(\varepsilon_{11} - \varepsilon_{22})^{2} + (\varepsilon_{22} - \varepsilon_{33})^{2} + (\varepsilon_{33} - \varepsilon_{11})^{2}] + \varepsilon_{12}^{2} + \varepsilon_{13}^{2} + \varepsilon_{23}^{2} \right\}^{1/2} .$$
(17)

The desired quantities used for comparison with laboratory tests are then obtained from Eqs. (9) - (12).

The results presented here have shown that

- (1) We can define strain paths for static testing of rock (and soil) samples that are more representative of actual field situations that those commonly used heretofore in constitutive modeling, and that
- (2) It is possible to reproduce these paths in laboratory tests.

APPENDIX I

General Relationships and Finite-Difference Calculations

The equation for momentum conservation in Eulerian coordinates is given by

$$-, \dot{\mathbf{v}} = \frac{\partial \sigma_{\mathbf{r}}}{\partial \mathbf{r}} + (g-1) \frac{\sigma_{\mathbf{r}} - \sigma_{i}}{\mathbf{r}} , \qquad (18)$$

where , is the material density, v is the radial particle velocity, σ_r and σ_θ are the radial and tangential stress components, and g is 1 (for plane flow), 2 (for cylindrical flow) or 3 (for spherical flow). A dot over a variable indicates time differentiation at a fixed material element and r is the Eulerian spatial coordinate. It is inconvenient to deal with Eulerian coordinates, hence we choose to express Eq. (18) in terms of Lagrangian coordinates representing the initial configuration. We define R as the initial radial coordinate of a material element whose current radial location is at r. Radial and transverse stress components in the initial configuration (Lagrangian) are denoted σ_R and σ_{Θ} . If the initial density is given by ρ_{Θ} , then mass conservation requires that

$$\rho_0 R^{g-1} dR = \rho r^{g-1} dr$$
 (19)

If the forces on a material element are to be the same in the two representations, then

$$R^{g-1} \sigma_R = r^{g-1} \sigma_r$$
 (20)

$$\sigma_{\Omega} dR^{g-1} = \sigma_{\theta} dr^{g-1} \qquad . \tag{21}$$

Now write Eq. (18) as

$$-\rho r^{g-1} dr \dot{v} = d(r^{g-1}\sigma_r) - \sigma_\theta dr^{g-1}$$
, (22)

keeping in mind that the differentials on the right-hand side are taken at constant time. Substitution of Eqs. (19) - (21) into Eq. (22) then gives

$$-\rho_{o}R^{g-1} dR \dot{v} = d(R^{g-1}\sigma_{R}) - \sigma_{\odot}dR^{g-1}$$
, (23)

or

$$-\rho_{0}\dot{\mathbf{v}} = \frac{\partial\sigma_{R}}{\partial R} + (g-1)\frac{\sigma_{R} - \sigma_{\Theta}}{R} , \qquad (24)$$

in Lagrangrian coordinates.

In order to use Eq. (24) in a finite-difference solution, an artificial viscous stress q is included. The following equations, with the addition of a constitutive law, then form the basis of the numerical calculations:

$$\rho_{O}\dot{\mathbf{v}} = \frac{\partial \sigma_{R}}{\partial R} - (g-1) \frac{(\sigma_{R} - \sigma_{O})}{R} - \frac{\partial q}{\partial R}$$
 (25)

$$q = \rho_0 A^2 (\Delta R)^2 \left| \frac{\partial \mathbf{v}}{\partial R} \right|^2, \quad \frac{\partial \mathbf{v}}{\partial R} \le 0$$

$$= 0, \quad \frac{\partial \mathbf{v}}{\partial R} > 0$$
(26)

$$\dot{\varepsilon}_{R} = -\frac{\partial v}{\partial R}$$
 , $\dot{\varepsilon}_{\Theta} = -\frac{v}{R}$, (27)

where A is nondimensional constant on the order of unity, ΔR is the spatial increment in the finite-difference solution, and ε_R and ε_Θ are the radial and tangential strain rates in the initial configuration. A straight-forward centered difference scheme is used and Eqs. (25) - (27) are written in

finite-difference form as

$$\rho_{0} \frac{v_{j}^{i+\frac{1}{2}} - v_{j}^{i-\frac{1}{2}}}{\Delta t} = -\frac{(\sigma_{R})_{j+\frac{1}{2}}^{i} - (\sigma_{R})_{j-\frac{1}{2}}^{i}}{\Delta R}$$

$$(g-1) \frac{(\sigma_{R})_{j+\frac{1}{2}}^{i} + (\sigma_{R})_{j-\frac{1}{2}}^{i} - (\sigma_{\Theta})_{j+\frac{1}{2}}^{i} - (\sigma_{\Theta})_{j-\frac{1}{2}}^{i}}{2R_{i}}$$

$$-\frac{q_{j+\frac{1}{2}}^{i-\frac{1}{2}}-q_{j-\frac{1}{2}}^{i-\frac{1}{2}}}{\Delta R},$$
 (28)

$$(\dot{\varepsilon}_{R})_{j+\frac{1}{2}}^{j+\frac{1}{2}} = \frac{v_{j}^{j+\frac{1}{2}} - v_{j+\frac{1}{2}}^{j+\frac{1}{2}}}{\Delta R}, \qquad (29)$$

$$(\dot{\varepsilon}_{\Theta})_{\mathbf{j}+\frac{1}{2}}^{\mathbf{i}+\frac{1}{2}} = -\frac{v_{\mathbf{j}}^{\mathbf{j}+\frac{1}{2}} + v_{\mathbf{j}+\mathbf{1}}^{\mathbf{j}+\frac{1}{2}}}{2R_{\mathbf{j}+\frac{1}{2}}}$$
(30)

The stress rates ($\dot{\sigma}_R$ and $\dot{\sigma}_\Theta$) are obtained from $\dot{\epsilon}_R$ and $\dot{\epsilon}_\Theta$, and therefore the stresses and strains are calculated from

$$\chi_{j+\frac{1}{2}}^{i+1} = \chi_{j+\frac{1}{2}}^{i} + \chi_{j+\frac{1}{2}}^{i+\frac{1}{2}} \Delta t \qquad , \tag{31}$$

where X represents σ_R , σ_Θ , ϵ_R and ϵ_Θ .

The constitutive model used here is expressed in terms of the principal stress and strain components σ_i and ϵ_i (i = 1, 2 and 3) with the following identification:

g = 1 (Plane Flow):

$$\dot{\varepsilon}_1 = -\partial v/\partial R$$
 , $\dot{\varepsilon}_2 = \dot{\varepsilon}_3 = 0$
 $\sigma_1 = \sigma_R$, $\sigma_2 = \sigma_3 = \sigma_Z$

g = 2 (Cylindrical Flow)

$$\dot{\epsilon}_1 = -\partial v/\partial R$$
 , $\dot{\epsilon}_2 = -v/R$, $\dot{\epsilon}_3 = 0$ $\sigma_1 = \sigma_R$, $\sigma_2 = \sigma_\Theta$, $\sigma_3 = \sigma_Z$

g = 3 (Spherical Flow):

$$\dot{\epsilon}_1 = -\partial v/\partial R$$
 , $\dot{\epsilon}_2 = \dot{\epsilon}_3 = -v/R$ $\sigma_1 = \sigma_R$, $\sigma_2 = \sigma_3 = \sigma_\Theta$.

Let us define the volume strain $\epsilon_{\rm V}$, the mean stress p, the stress deviators s; and the second invariant of the stress tenser according to

$$\epsilon_{V} = \epsilon_{1} + \epsilon_{2} + \epsilon_{3}$$
 , (32)

$$p = (\sigma_1 + \sigma_2 + \sigma_3)/3 , (33)$$

$$s_{i} = \sigma_{i} - p \qquad , \tag{34}$$

$$J_2 = (s_1^2 + s_2^2 + s_3^2)/2 \qquad . \tag{35}$$

The elastic-plastic constitutive relation used here is then defined according to the following equations:

$$p = \hat{p}(\epsilon_{V}) , \qquad (36)$$

$$\dot{s}_{i} = 2\mu(\dot{\epsilon}_{i} - \dot{\epsilon}_{v}/3) - 2\mu\xi\frac{s_{i}}{\sqrt{J_{2}}} \qquad (37)$$

The variable ξ is determined by the condition that the stress state must remain on the failure surface, defined by

$$\sqrt{J_2} = f(p) \qquad , \tag{38}$$

when a material element is undergoing plastic deformation.

From Eq. (35) we find that

$$2\sqrt{J_2} / \sqrt{J_2} = s_i s_i$$
 (Summation) (39)

and

$$\sqrt{J_2} = (\mu/\sqrt{J_2}) s_i \dot{\epsilon}_i - 2\mu\xi = f'(p)\dot{p}$$
 (40)

Therefore, the variable ξ in Eq. (37) is given by

$$2\mu\xi = (\mu/\sqrt{J_2}) s_{i}\dot{\epsilon}_{i} - f'(p)p$$
 , (41)

or, in terms of $\boldsymbol{\sigma}_{\boldsymbol{i}}$ and $\boldsymbol{p}_{\boldsymbol{i}}$ as

$$2\mu\xi = (\mu/\sqrt{J_2})(\sigma_{\dot{1}}\dot{\epsilon}_{\dot{1}} - p\dot{\epsilon}_{\dot{V}}) - f'(p)\dot{p} . \qquad (42)$$

If it is desired to include effects of fluid saturation defined by nonzero pore pressure p_p , σ_i is replaced by the effective stress components $\langle \sigma_i \rangle \equiv \sigma_i - nP_p \ (0 < n < 1)$ in the elasticity relationship and by $\sigma_i^* \equiv \sigma_i - P_p$ in the failure surface relationship:

$$\langle p \rangle = p - nP_p = \hat{p}(\epsilon_v)$$
 (43)

$$\langle s_{i} \rangle = s_{i} = 2\mu(\dot{\epsilon}_{i} - \dot{\epsilon}_{v}/3) - 2\mu\xi\frac{s_{i}}{\sqrt{J_{2}}},$$
 (44)

$$2\mu\xi = (\mu/\sqrt{J_2})(\sigma_i\dot{\epsilon}_i - p\dot{\epsilon}_v) - f'(p^*)(1-m)\dot{p}$$
, (45)

where

$$m \equiv \frac{dP_p}{dp} \qquad . \tag{46}$$

The function f(p) is taken to be of the form

$$f(p) = S_0 + \Delta S(1 - e^{-p/a})$$
 (47)

<u>Analytical Determination of Elastic Stress and Strain Paths for a Spherical Explosion</u>

If u(r,t) is the radial displacement, the spherical wave equation for purely elastic deformation can be written as

$$\partial^2 u / \partial t^2 = c^2 [\partial^2 u / \partial r^2 + (2/r) \partial u / \partial r - (2/r^2) u]$$
, (48)

where r is the radial coordinate, t is the time and c is the longitudinal elastic wave speed. This expression takes a simpler form if it is written in terms of a displacement potential Ψ such that

$$u(r,t) = c^2 \partial/\partial r (\Psi/r) \qquad . \tag{49}$$

In this case

$$a^2 \Psi / at^2 = c^2 a^2 \Psi / ar^2$$
 (50)

whose solution for outgoing waves is given by the familiar expression

$$\Psi = \Psi(t - \frac{r - r_0}{c}) \qquad . \tag{51}$$

The displacement, strain components and stress components can be expressed in terms of Ψ and its derivatives Ψ' and Ψ'' according to

$$u(r,t) = -(c/r)\psi' - (c/r)^2 \psi$$
, (52)

$$-\varepsilon_a = \partial u/\partial r = (1/r)\Psi'' + (2c/r^2)\Psi' + (2c^2/r^3)\Psi$$
 , (53)

$$-\varepsilon_{t} = u/r = -(c/r^{2})\Psi' - (c^{2}/r^{3})\Psi$$
 , (54)

$$-\sigma_{a} = (1/r) [(\lambda + 2\mu) \Psi'' + (4\mu c/r) \Psi' + (4\mu c^{2}/r^{2}) \Psi] , \qquad (55)$$

$$-\sigma_{t} = (1/r) [\lambda \Psi'' - (2\mu c/r)\Psi' - (2\mu c^{2}/r^{2})\Psi]$$
 , (56)

where λ and μ are the Lame constants. The sign convention used throughout this work is that stresses and strains are positive in compression. For a pressure history at $r = r_0$ given by

$$\sigma_{\mathbf{r}}(\mathbf{r}_{0}, \mathbf{t}) = 0 , \qquad \mathbf{t} < 0$$

$$\sigma_{\mathbf{r}}(\mathbf{r}_{0}, \mathbf{t}) = \rho_{0} e^{-\alpha \mathbf{t}} , \qquad \mathbf{t} \ge 0$$

$$(57)$$

The function Ψ must satisfy the following ordinary differential equation:

$$(\lambda + 2\mu) \Psi''(t) + (4\mu c/r_0) \Psi'(t) + (4\mu c^2/r_0^2) \Psi(t) = (58)$$

$$-r_0 p_0 e^{-\alpha t} ,$$

subject to the conditions, from Eqs. (52) and (58), that jumps in Ψ and Ψ' at t = 0 obey the following relationships:

$$(\lambda + 2\mu) [\Psi'] + (4\mu c/r_0) [\Psi] = 0$$
,
 $[\Psi'] + (c/r_0) [\Psi] = 0$, (59)

where [] indicates the jump in the function, i.e., [f] = $f(0^+)$ - $f(0^-)$. Equations (59) thus require that Ψ and Ψ' each be continuous at t = 0 as long as $\lambda \neq 2\mu$. Hence, a solution to Eq. (58) can be written as

$$\Psi(t) = e^{-\beta_2 t} \left(M \cos \beta_1 t + N \sin \beta_1 t \right) + \Psi_0 e^{-\alpha t} , \qquad (60)$$

where

$$M = -\Psi_0 = \frac{r_0 p_0}{\alpha^2 (\lambda + 2\mu) - 4\mu c\alpha/r_0 + 4\mu c^2/r_0^2},$$
 (61)

$$N = \frac{\alpha r_0 (\lambda + 2\mu) - 2\mu c}{2c \sqrt{\mu(\lambda + \mu)}} \Psi_0 \qquad , \tag{62}$$

$$\beta_1 = \frac{2c \sqrt{\mu(\lambda + \mu)}}{r_0 (\lambda + 2\mu)} , \qquad (63)$$

$$\beta_2 = \frac{2\mu c}{r_0 (\lambda + 2\mu)} \qquad (64)$$

In the case of an elastic fluid μ = 0 and the displacement potential and its first two derivatives become

$$\Psi = \frac{r_0 p_0}{\lambda \alpha^2} (1 - e^{-\alpha t} - \alpha t) \qquad , \tag{65}$$

$$\Psi' = \frac{r_0 p_0}{\lambda \alpha} \left(e^{-\alpha t} - 1 \right) \qquad , \tag{66}$$

$$\Psi' = -\frac{r_0 p_0}{\lambda} e^{-\alpha t} \qquad . \tag{67}$$

If $\alpha = 0$ (i.e., the cavity pressure remains constant at p_0) in the case of

a fluid, the displacement potential and its first two derivatives become

$$\Psi = -\frac{r_0 p_0}{2\lambda} t^2 \qquad , \tag{68}$$

$$\Psi' = -\frac{r_0 p_0}{\lambda} t \qquad , \tag{69}$$

$$\Psi'' = -\frac{r_0 p_0}{\lambda} \qquad . \tag{70}$$

In the special case of spherical wave propagation we can make the identification that $L = \sigma_a - \sigma_t$ and $p_c = \sigma_t$, in which case the stress and strain paths can be written parametrically as

$$L = -(2\mu/r)[\Psi'' + (3c/r)\Psi' + (3c^2/r^2)\Psi] , \qquad (71)$$

$$P_c = -(1/r)[\lambda \Psi'' - (2\mu c/r)\Psi' - (2\mu c^2/r^2)\Psi]$$
 , (72)

$$\varepsilon_a = -(1/r)[\Psi'' + (2c/r)\Psi' + (2c^2/r^2)\Psi]$$
, (73)

$$\varepsilon_{\dagger} = (c/r^2)[\Psi' + (c/r)\Psi] \qquad . \tag{74}$$

Equations (71) to (74) in the case of spherical elastic waves are the analytical counterparts of Eqs. (9) to (12) for numerical solutions. Comparison of strain and stress paths calculated by the two methods is shown in Figure 8 for $1/\alpha = 1$ msec, $R/R_0 = 3$, K = 95 kbar, C = 3 km/sec, and $P_0 = 2.0$ gm/cm³. It can be seen that the numerical solution gives a good approximation of the strain and stress paths except for the peak values associated with the main compressive fronts. This is a result of the viscous stresses that are included in the finite-difference solution to

damp out numerical oscillations, and has no significance with regard to the conclusions reached in this report.

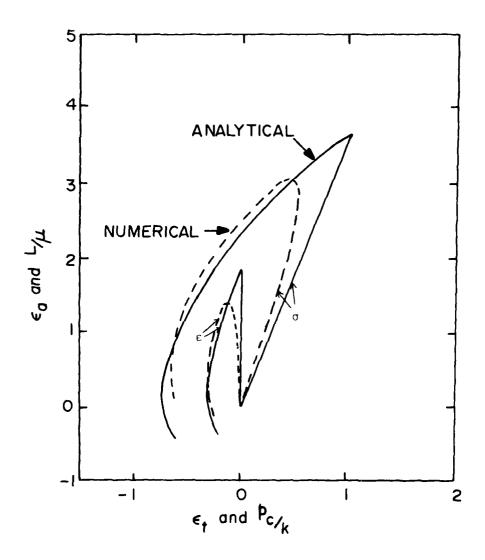


Figure 8. Comparison of strain and stress paths determined numerically and analytically for spherical wave propagation in an elastic medium.

APPENDIX II

EXPERIMENTAL TECHNIQUE

Specimen Preparation

Specimens were prepared from Kayenta sandstone, Mixed Company Site. Cylindrical samples 3.81 centimeters long by 1.91 centimeters diameter were used thus maintaining a length to diameter ratio of 2 to 1. Specimen ends were ground parallel to within \pm .001 centimeters. Specimens were air dried with weight, length and diameters being recorded for each sample for use in determining sample density and strains. Samples were prepared for testing by first wrapping them in urethane plastic (.025 cm thick) with hardened steel endcaps attached at each end using stainless steel lock wire.

Stress and Strain Determination

Stress and strain transducers were placed within the pressure vessel. Confining pressure was measured using a calibrated 350-ohm manganin pressure sensitive coil accurate to \pm .003 kbars. Jacketed samples were placed and centered on the load cell when in the pressure vessel. The load cell was accurate to \pm .005 kbars. Axial and lateral strain transducers were of the cantilever type using strain gauges in a wheatstone bridge configuration to obtain voltage output. The axial cantilevers measured total axial displacement and were calibrated to be accurate to \pm .003 percent strain. Lateral strain cantilevers were positioned at mid-sample and sampled strains at 90 degree intervals. Diametrically opposed arms were calibrated for lateral strain. The lateral strains were averages with a resulting accuracy of \pm .006

percent. Figure 9 shows a schematic of the transducers when inside the pressure vessel. Further discussion on transducer design may be obtained in Terra Tek report TR 75-29.

Testing Procedures

Seven samples were first tested triaxially to failure to generate the triaxial failure envelope for the material while eight samples were tested following the three strain paths. Triaxial testing commenced by first hydrostatically loading the samples to the desired confining pressure with subsequent axial loading to failure, stresses and strains being recorded during all phases of loading. A strain rate of about $10^{-4}~{\rm sec}^{-1}$ was used during loading.

Uniaxial-strain loading was used when following a specified strain path. Axial load and confining pressure were applied such that zero lateral strain was maintained. When following strain path I, II or III during unloading, i.e., constant-axial-strain and uniaxial-strain unloading, constant axial strain unloading and constant volume strain unloading, respectively, the confining pressure and axial load were adjusted to maintain the desired strain state.

Data Acquisition and Analysis

Both x-y recorders and a PDP Lab 11 computer were used for data acquisition. The x-y recorders were used primarily for instantaneous feedback during testing while the PDP Lab 11 computer data was used for analysis of pressure effects, endcap effects and generation of stress and strain load-unload curves. Tables I, II and III presented in the text are a result of the computer analysis.

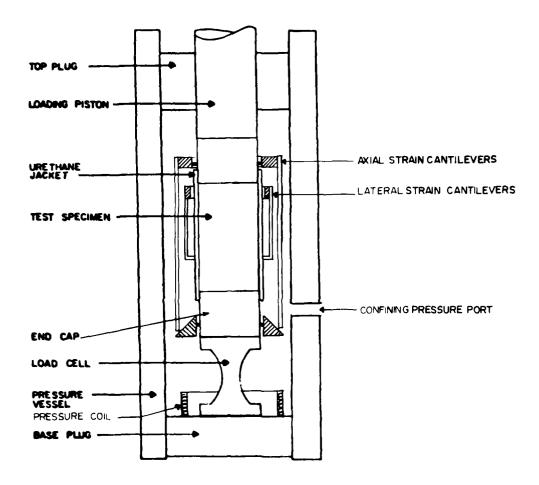


Figure 9. Pressure vessel schematic showing the sample and stress and strain transducers.

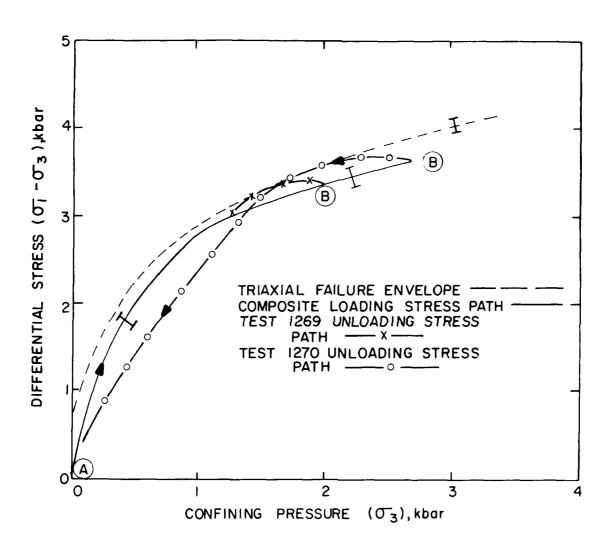


Figure 9a. Stress path followed during strain path III testing.

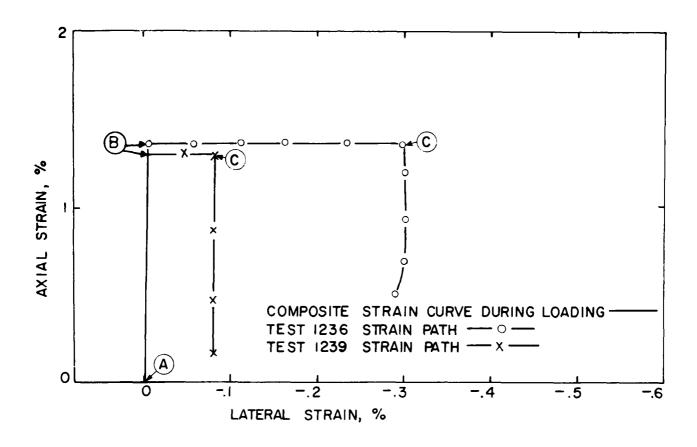


Figure 9b. Strain path followed during path I testing.

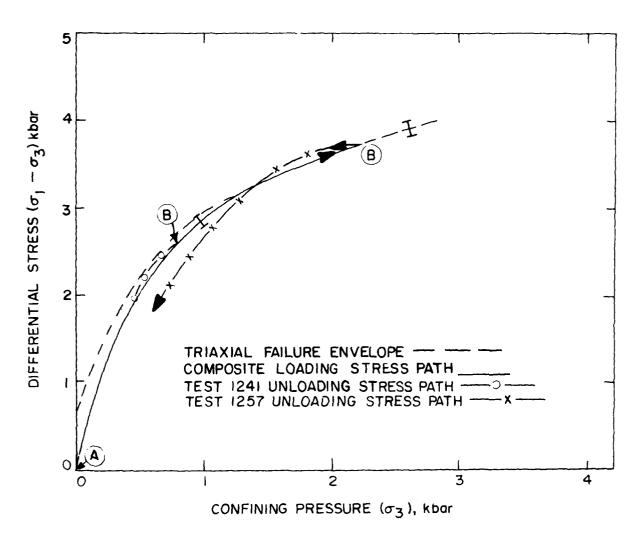


Figure 9c. Stress path followed during strain path II testing.

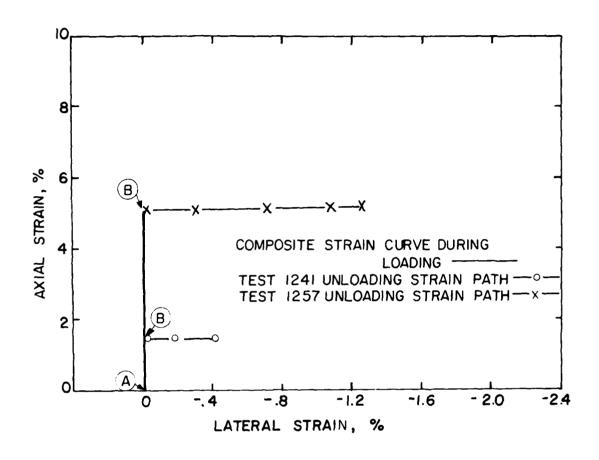


Figure 9d. Strain path followed during path II testing.

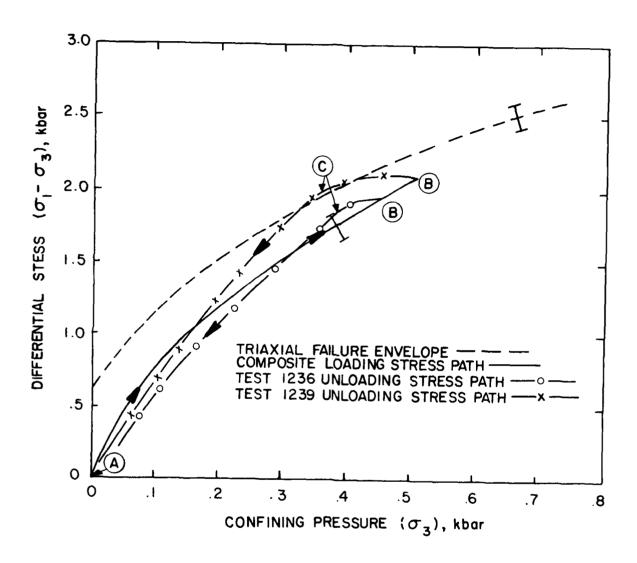


Figure 9e. Stress path followed during strain path I testing.

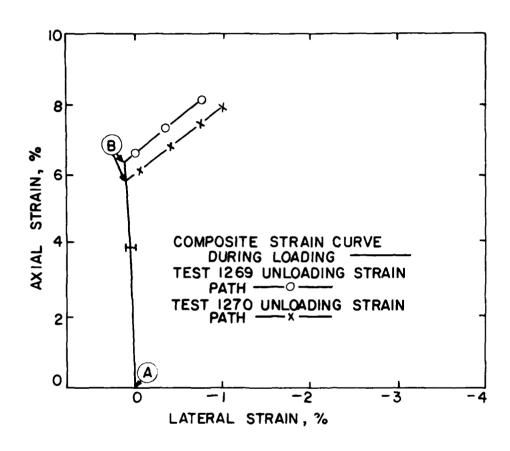


Figure 9f. Strain path followed during path III testing.

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